Final Examination, 12 Dec 2002 (Solutions) SM311O (Fall 2002)

1. (a) Let $\mathbf{v} = \langle x^2 z, \sin(ay), z\sqrt{x} \rangle$ where a is a constant. Determine a so that the divergence of \mathbf{v} vanishes at the point P = (4, 0, 1).

Solution: div $\mathbf{v} = 2xz + a\cos(ay) + \sqrt{x}$. At P the divergence has the value 10 + a, which vanishes when a = -10.

- (b) Let $\mathbf{v} = \langle y^2 x, y x^2, 0 \rangle$. Find the curl of \mathbf{v} . Is this flow irrotational anywhere?
- (c) Prove the identity $\nabla \times \nabla \phi = \mathbf{0}$ if ϕ is an arbitrary function of x, y, and z.

Solution: $\nabla \times \nabla \phi = \nabla \times \langle \phi_x, \phi_y, \phi_z \rangle = \langle \phi_{zy} - \phi_{yz}, \phi_{xz} - \phi_{zx}, \phi_{yx} - \phi_{xy} \rangle = \langle 0, 0, 0 \rangle$ because the order of differentiation does not matter for smooth (at least twice differentiable) functions.

- 2. Verify by direct differentiation if
 - (a) $u(z) = e^{2z} \cos 2z$ is a solution of u'' + au' + bu = 0 for any pair (a, b).

Solutions: With $u = e^{2z} \cos 2z$ the differential operator u'' + au' + bu takes the form

$$e^{2z} ((2a+b)\cos(2z) - 2(4+a)\sin(2z)).$$

- . This expression vanishes for ALL z if 2a + b = 0 and a + 4 = 0, or if a = -4 and b = 8.
- (b) $u(x,y) = \sin 3x \cos 4y$ is an eigenfunction of the Laplace operator $-\frac{\partial^2}{\partial x^2} \frac{\partial^2}{\partial y^2}$. What is the eigenvalue?

Solution: A function u is an eigenfunction of $-\frac{\partial^2}{\partial x^2} - \frac{\partial^2}{\partial y^2}$ with eigenvalue λ if $-\frac{\partial^2 u}{\partial x^2} - \frac{\partial^2 u}{\partial y^2} = \lambda u$. Let $u = \sin 3x \cos 4y$. Then $-\frac{\partial^2 u}{\partial x^2} - \frac{\partial^2 u}{\partial y^2} = 25 \sin 3x \cos 4y = 25u$. So u is an eigenfunction with eigenvalue 25.

3. (a) Give a parametrization for the plane that passes through the points (1,1,0), (0,2,2), and (3,0,3).

Solution: First we use the three points on the plane and find two vectors that are parallel with the plane: (1,1,0) and (0,2,2) give $\mathbf{r}_1 = \langle -1,1,2\rangle$; (0,2,2) and (3,0,3) give $\mathbf{r}_2 = \langle 3,-2,1\rangle$. Next we form the cross product of \mathbf{r}_1 and \mathbf{r}_2 to get \mathbf{n} a normal vector to the plane:

$$\mathbf{n} = \mathbf{r}_1 \times \mathbf{r}_2 = \langle 5, 7, -1 \rangle.$$

Let (x, y, z) be any point on the plane. Considering that (1, 1, 0) is also on the plane, the vector $\mathbf{r}_3 = \langle x - 1, y - 1, z \rangle$ is parallel with the plane. Then $\mathbf{n} \cdot \mathbf{r}_3 = 0$ i.e. 5x + 7y - z = 12.

(b) Find a unit normal vector to the surface of the upper hemisphere of the Earth at the point whose longitude and latitude are 45 and 60 degrees, respectively.

Solution: The upper hemisphere is parametrized as

$$\mathbf{r}(u, v) = R\langle \cos u \cos v, \sin u \cos v, \sin v \rangle$$

where R is the Earth's radius and u and v are longitude and latitude. Then $\mathbf{n} = \mathbf{r}_u \times \mathbf{r}_v$ is normal to the hemisphere. Now $\mathbf{r}_u \times \mathbf{r}_v = \langle R^2 \cos(u) \cos(v)^2, R^2 \cos(v)^2 \sin(u), R^2 \cos(v) \sin(v) \rangle$. The

magnitude of this vector is $R^2 \cos v$. Dividing **n** by its magnitude yields the desired unit vector $\mathbf{N} = \langle \cos(u) \cos(v), \cos(v) \sin(u), \sin(v) \rangle$. Finally evaluating at $u = \frac{\pi}{4}$ and $v = \frac{\pi}{3}$ gives $\mathbf{N} = \langle \frac{1}{2\sqrt{2}}, \frac{1}{2\sqrt{2}}, \frac{\sqrt{3}}{2} \rangle$.

- 4. (a) The function $\phi(x,y) = ax^2y^2 by^2 + ax by$ is the potential for a velocity vector field \mathbf{v} . Determine all values of a and b so that the velocity of the particle located at (2,-1) is $\langle 1,2\rangle$. Solution: $\mathbf{v} = \nabla \phi = \langle 2axy^2 + a, 2ax^2y 2by b \rangle$. Evaluating \mathbf{v} at (2,-1) yields $\langle 5a, -8a + b \rangle$ which equals $\langle 1,2 \rangle$ if $a = \frac{1}{2}$ and $b = \frac{18}{5}$.
 - (b) The function $\psi(x,y) = ax^2 + xy + by^2$ is the stream function of a velocity field \mathbf{v} . Find all a and b so that the velocity of the particle located at (1,-2) has magnitude $\frac{1}{2}$. **Solution**: $\mathbf{v} = \langle \frac{\partial \psi}{\partial y}, -\frac{\partial \psi}{\partial x} \rangle = \langle x+2by, -2ax-y \rangle$. Evaluating this vector at (1,-2) and setting its magnitude equal to $\frac{1}{2}$ yields the equation $(2-2a)^2 + (1-4b)^2 = \frac{1}{4}$ for the set of all (a,b)s.
- 5. (a) Consider the velocity field $\mathbf{v} = (x^2z x)\mathbf{k}$. Determine the flux of this fluid through the following two surfaces:
 - i. a disk of radius 1 in the xy-plane and centered at the origin. **Solution**: First parametrize the disk: $\mathbf{r}(u,v) = \langle u\cos v, u\sin v, 0 \rangle$. Next compute $\mathbf{r}_u \times \mathbf{r}_v = \langle 0,0,u \rangle$. Then $\int \int_S \mathbf{v} \cdot d\mathbf{r} = \int_0^{2\pi} \int_0^1 \mathbf{r}|_S \cdot \mathbf{r}_u \times \mathbf{r}_v \, du \, dv = \int_0^{2\pi} \int_0^1 \langle 0,0,-u\cos v \rangle \cdot \langle 0,0,u \rangle \, du \, dv = 0$ so as much fluid is passing through S from below to above it as in the opposite direction.
 - ii. a disk of radius 1 in the plane z=3 and centered at the origin. **Solution**: First parametrize the disk: $\mathbf{r}(u,v)=\langle u\cos v,u\sin v,3\rangle$. Next compute $\mathbf{r}_u\times\mathbf{r}_v=\langle 0,0,u\rangle$. Then $\int\int_S\mathbf{v}\cdot d\mathbf{r}=\int_0^{2\pi}\int_0^1\mathbf{r}|_S\cdot\mathbf{r}_u\times\mathbf{r}_v\,dudv=\int_0^{2\pi}\int_0^1\langle 0,0,-u\cos v+3u^2\cos^2v\rangle\cdot\langle 0,0,u\rangle\,dudv=\frac{3\pi}{4}$.
 - (b) Use the Stokes Theorem or compute the appropriate surface integral to determine the flux of vorticity of $\mathbf{v} = x^2 \mathbf{k}$ through the surface of the upper hemisphere of a sphere of radius 2 centered at the origin.

Solution 1) Direct computation: $\omega = \nabla \times \mathbf{v} = \langle 0, -2x, 0 \rangle$. $\mathbf{r}(u, v) = \langle 2 \cos u \cos v, 2 \sin u \cos v, 2 \sin v \rangle$. $\mathbf{r}_u \times \mathbf{r}_v = \langle 4 \cos u \cos^2 v, 4 \cos^2 v \sin u, 2 \sin 2v \rangle$. Then $\int \int_S \omega \cdot d\mathbf{A} = \int_0^{2\pi} \int_0^{\frac{\pi}{2}} \langle 0, -4 \cos u \cos v, 0 \rangle \cdot \langle 4 \cos u \cos^2 v, 4 \cos^2 v \sin u, 2 \sin 2v \rangle dv du = 0$.

2) Using the Stokes Theorem: We need to compute the line integral $\int_C \mathbf{v} \cdot d\mathbf{r}$. $\mathbf{r} = \langle 2\cos t, 2\sin t, 0 \rangle$. But $\mathbf{v}|_C = \langle 0, 0, 4\cos^2 t \rangle$ is orthogonal to $\mathbf{r}'(t) = \langle -2\sin t, 2\cos t, 0 \rangle$. So the line integral will vanish.

6. Consider the following wave equation initial-boundary value problem:

$$u_{tt} = 4u_{xx},$$
 $u(0,t) = u(3,t) = 0,$ $u(x,0) = x(3-x),$ $u_t(x,0) = 0.$

(a) Use separation of variables and find the solution to this problem. Clearly indicate the process of separation of variables and the Fourier Series method used in obtaining this solution.

Solution: From separation of variables we get that

$$u(x,t) = \sum_{n=1}^{\infty} (a_n \cos \frac{2n\pi t}{3} + b_n \sin \frac{2n\pi t}{3}) \sin \frac{n\pi x}{3}.$$

But $u_t(x,0) = 0$ which implies $b_n = 0$ for all n. u(x,0) = x(3-x) so

$$a_n = \frac{(x(3-x), \sin\frac{n\pi x}{3})}{(\sin\frac{n\pi x}{3}, \sin\frac{n\pi x}{3})} = \frac{2}{3} \, \left(\frac{54}{n^3 \, \pi^3} - \frac{54 \, \cos(n \, \pi)}{n^3 \, \pi^3}\right).$$

So

$$u(x,t) = \frac{2}{3} \sum_{n=1}^{\infty} \left(\frac{54}{n^3 \pi^3} - \frac{54 \cos(n \pi)}{n^3 \pi^3} \right) \cos \frac{2n\pi t}{3} \sin \frac{n\pi x}{3}.$$

(b) Use the first nonzero term of the above solution and estimate u(3/2, 3/4).

Solution: With n = 1 the solution takes the form

$$u(x,t) = \frac{72}{\pi^2} \cos \frac{2\pi t}{3} \sin \frac{\pi x}{3}.$$

So u(3/2, 3/4) = 0.

(c) Use the first nonzero term of the above solution and estimate how long it takes for the wave to go though one complete vibration.

Solution: The period of $\cos \frac{2\pi t}{3}$ is $2\pi/(2\pi/3) = 3$.

7. Let $\mathbf{v} = \langle x^2 + y^2, 2xy \rangle$ be the velocity field of a fluid. Compute the acceleration \mathbf{a} of this flow. Does \mathbf{a} have a potential p? If yes, find it.

Solution: $\mathbf{a} = \mathbf{v} \cdot \nabla v = \langle 2(x^3 + 3xy^2), 2(3x^2y + y^3) \rangle$. This vector has zero curl so there is a function p such that $\nabla p = \langle 2(x^3 + 3xy^2), 2(3x^2y + y^3) \rangle$. From

$$\frac{\partial p}{\partial x} = 2(x^3 + 3xy^2), \quad \frac{\partial p}{\partial y} = 2(3x^2y + y^3),$$

we get that $p = \frac{x^4}{2} + 3x^2y^2 + \frac{y^4}{2}$.

- 8. Let Ω stand for the angular velocity of our planet.
 - (a) Noting that our planet rotates once every 24 hours, compute Ω where $\Omega = \langle 0, 0, \Omega \rangle$. What are the units of Ω ?

Solution: $\Omega = \frac{2\pi \, \text{radians}}{(24 \, \text{hours})(60 \, \text{minutes/hour})(60 \, \text{seconds/minute})} = 0.00007 \, \text{rad/s}.$

(b) Use this value of Ω and estimate the values in the centripetal acceleration $\Omega \times (\Omega \times \mathbf{r})$ where \mathbf{r} is the position vector to a typical point on the surface of the Earth. Assume that the radius of the Earth is 6000 kilometers.

Solution: Let P = (x, y, z) be a point on the planet. Note that $\mathbf{\Omega} = \langle 0, 0, \Omega \rangle$. Then $\mathbf{r} = \langle x, y, z \rangle$ and $\mathbf{\Omega} \times (\mathbf{\Omega} \times \mathbf{r}) = -\Omega^2 \langle x, y, 0 \rangle$. x or y take their largest values at the equator which could be as large as 6000 kilometers. So the nonzero components of the centripetal acceleration could be as large as $0.00007^2 \, (\text{rad/s})^2 \times 6000000 \, \text{meters} = 0.0294 \, \text{m/s}^2$, considerably smaller than $9.8 \, \text{m/s}^2$ from the acceleration of gravity.

9. Consider an incompressible fluid occupying the basin

$$D = \{(x, y, z) | 0 \le z \le H\}.$$

Let $\mathbf{v} = \langle v_1, v_2, v_3 \rangle$ be the velocity field of a motion generated in D. Suppose that we have been able to determine that

$$v_1(x, y, z) = 3x^2y^2 - x, \quad v_2(x, y, z) = yz,$$

but have only succeeded in measuring v_3 at the bottom of the basin and that this value is

$$v_3(x, y, 0) = x + y.$$

Determine v_3 everywhere in D. (Hint: What does incompressibility mean **mathematically**?)

Solution: From the equation of incompressibility we have

$$\frac{\partial v_1}{\partial x} + \frac{\partial v_2}{\partial y} + \frac{\partial v_3}{\partial z} = 0$$

or

$$\frac{\partial v_3}{\partial z} = -\frac{\partial v_1}{\partial x} - \frac{\partial v_2}{\partial y}.$$

Substituting the values of v_1 and v_2 in the above relation yields

$$\frac{\partial v_3}{\partial z} = -2xy^2.$$

Integrating this result with respect to z from 0 to z and using the value of v_3 at z=0 yields

$$v_3 = -2xy^2z + x + y.$$

Solution: Since the flow is incompressible,

10. A flow is called geostrophic if the velocity $\mathbf{v} = \langle u(x,y), v(x,y) \rangle$ and the pressure gradient ∇p are related by

$$(*) \qquad -fv = -rac{1}{
ho}rac{\partial p}{\partial x}, \quad fu = -rac{1}{
ho}rac{\partial p}{\partial y},$$

where ρ , a constant, is the density of the fluid, and f is the coriolis parameter.

(a) Assuming that f is constant, prove that the divergence of \mathbf{v} must vanish.

Solution: From the equations of motion we have

$$u = -\frac{1}{\rho f} \frac{\partial p}{\partial y}, \quad v = \frac{1}{\rho f} \frac{\partial p}{\partial x}.$$

Now div $\mathbf{v} = \frac{\partial v_1}{\partial x} + \frac{\partial v_2}{\partial y}$ which is equal to

$$-\frac{1}{\rho f}\frac{\partial^2 p}{\partial x \partial y} + \frac{1}{\rho f}\frac{\partial^2 p}{\partial y \partial x} = 0$$

- (b) Prove that the particle paths of a geostrophic flow and its isobars coincide.
 - **Solution**: Note that $\mathbf{v} \cdot \nabla p = -\frac{1}{\rho f} \frac{\partial p}{\partial y} \frac{\partial p}{\partial x} + \frac{1}{\rho f} \frac{\partial p}{\partial x} \frac{\partial p}{\partial y} = 0$. So \mathbf{v} and ∇p are orthogonal. Since ∇p is orthogonal to isobars, and since \mathbf{v} is tangential to particle paths, particle paths and isobars coincide.
- (c) Consider a high pressure field in a geostrophic flow in the northern hemisphere (where f > 0). By appealing to the equations in (*) explain whether this high pressure field results in a clockwise or a counterclockwise motion.

Solution: Without loss of generality, assume that the high pressure occurs at the origin of the coordinate system. Let P be a point in the first quadrant. Then ∇p at P points toward the origin because 0 is a maximum of p. Then $\frac{\partial p}{\partial x} \leq 0$ and $\frac{\partial p}{\partial y} \leq 0$ at P (draw a picture to convince yourself of this). Going back to the geostrophic equations, $u \geq 0$ and $v \leq 0$ at P which indicates that the motion is clockwise.